

Reduced-model (SOLT) simulations of an EDA H-mode shot at C-Mod[†]

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presented at the TTF Meeting, Apr. 6 - 9, 2011, San Diego

[†]Work supported by USDOE Grant No. DE-FG02-97ER54392, USDOE Contract No. DE-AC02-09CH11466, USDOE Cooperative Agreement No. DE-FC02-99ER54512 and PPPL Subcontract No. S009625-F.

1) SOLT recovers the observed **heat-flux width scaling** of
EDA H-modes at C-Mod

2) SOLT's quasi-coherent mode (QCM)

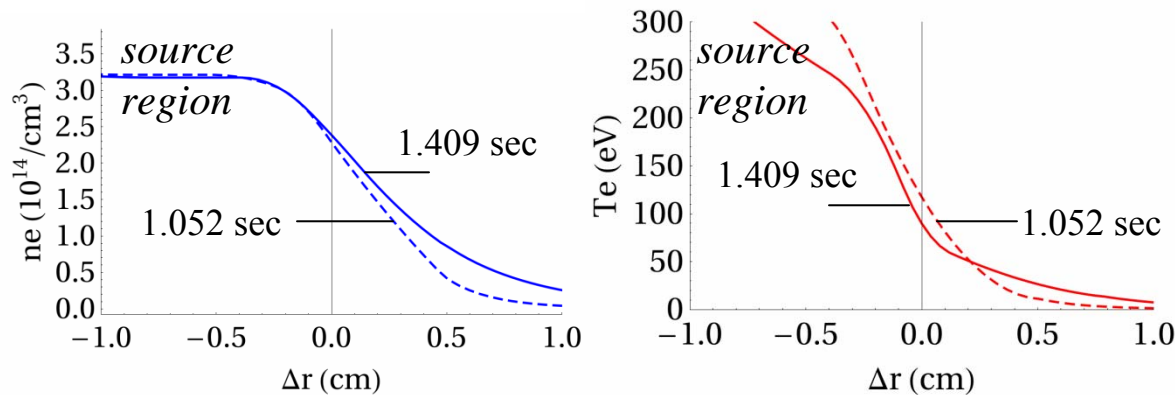
The Model

Scrape-Off-Layer Turbulence (SOLT) model simulations

- 2D \perp to B in OM
- electrostatic fluid model, reduced from Braginskii
- sheath physics (closure relations)
- Turbulent, O(1) fluctuations (n_e , T_e , ϕ)
- mean poloidal flows (p_y) from momentum conservation,
with sheath physics and viscosity
- blobs, EDWs, profile modification

Input to SOLT from Experiment

- C-Mod profiles (n_e , T_e) for EDA H-mode shot #1100303018 at two time slices :



- SOLT profiles are damped to these for $\Delta r < 0$, otherwise they evolve by (self-consistent) SOLT dynamics.

- We add an adjustable mean flow (ZF): $\langle v_y \rangle_y \equiv \bar{v}_y(\Delta r, t)$ based on ion pressure balance.

- In the SOL ($\Delta r > 0$), the flow evolves by momentum conservation and sheath physics.
- On the core-side ($\Delta r < 0$), the flow is damped to a reference, \bar{v}_{y0} , derived from the C-Mod profiles.

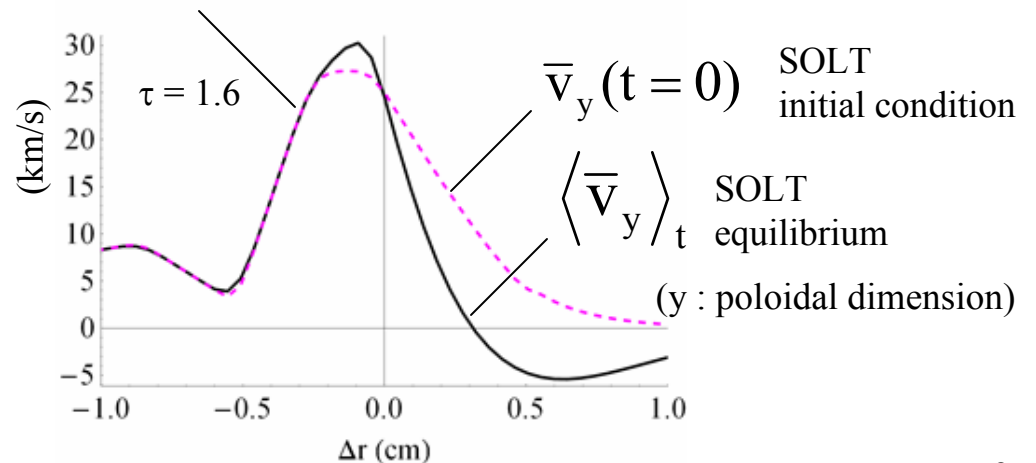
Reference Flow :

$$e Z n_i E_r - \partial_r (n_i T_i) = 0$$

$\Rightarrow E \times B$ drift :

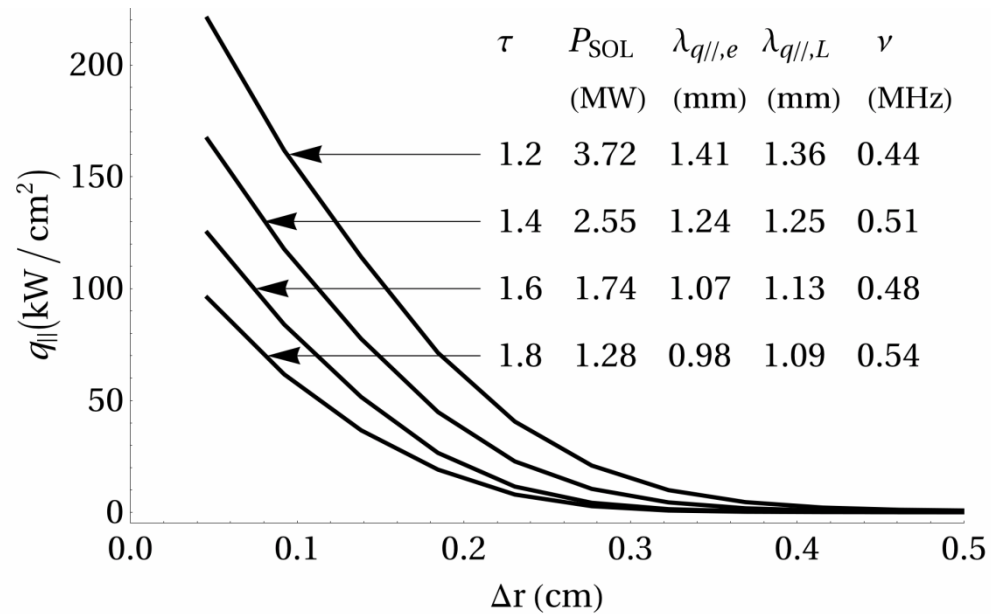
$$\bar{v}_{y0} = -\tau \cdot \partial_r (n_e T_e) / n_e, \tau \sim T_i / T_e$$

τ controls the turbulence



τ-Scan Results from SOLT

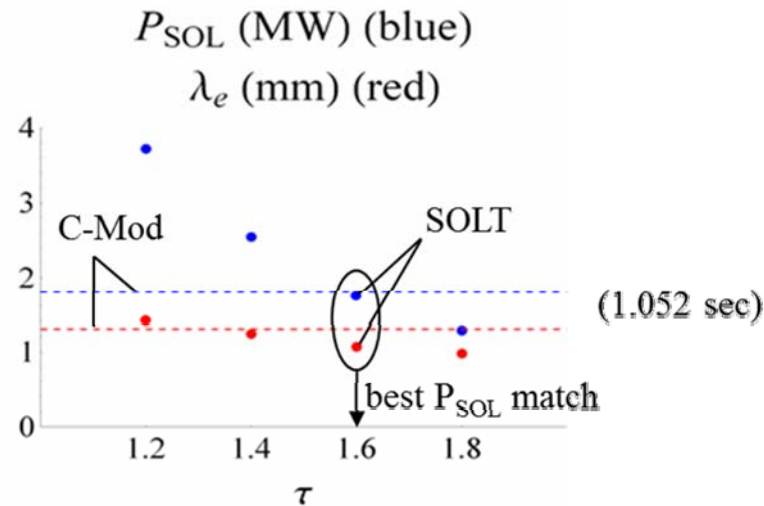
Change τ , the amplitude of \bar{V}_{y0} , in SOLT to scan P_{SOL} and q_{\parallel} ($v_{\parallel}/n_e T_e$).



$\lambda_{q_{\parallel},e}$: exponential fit; $\lambda_{q_{\parallel},L}$: Loarte length

ν : location of the peak in the density fluctuation spectrum at $\Delta r = 0.46$ mm.

τ-Scan Results from SOLT (cont.)



For the best P_{SOL} match at each time-slice :

t (s)	$P_{\text{C-Mod}}$	P_{SOLT}	$T_{\text{C-Mod}}$	T_{SOLT}	$\lambda_{e \text{ C-Mod}}$	$\lambda_{e \text{ SOLT}}$
1.052	1.79	1.74	117	63	1.31	1.07
1.409	1.39	1.32	86	48	2.27	1.57

SOL width (λ_e) decreases with increasing power (and T_{sep})
in both experiment and simulation.

Parallel Heat Flux is Limited by Collisions in the near-SOL

parallel heat flux regimes

- flux - limited :

$$q_{FL} = C_{FL} n_e v_e T_e$$

- sheath - limited :

$$q_{SL} = s_E n_e c_s T_e \exp[e(\Phi_B - \Phi)/T_e]$$

- collision - limited :

$$q_{CL} = 3.2 n_e c_s T_e / \Lambda$$

- $1/q_{//} = 1/q_{FL} + 1/q_{SL} + 1/q_{CL}$

The bottleneck sets the regime.

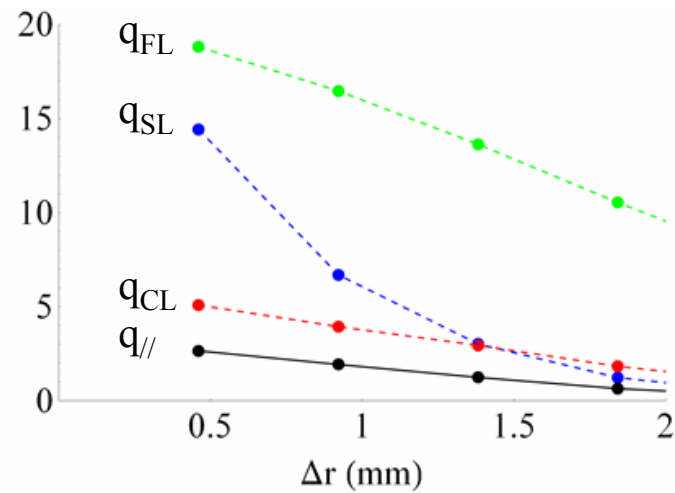
$$\nabla \cdot q = 0$$

⇓

$$\lambda_e \sim L_{//} q_{\perp} / q_{//}$$

$$q_{CL} \sim T_e^{7/2} \Rightarrow \text{smaller SOL widths } (\lambda_e) \text{ at higher } T_e$$

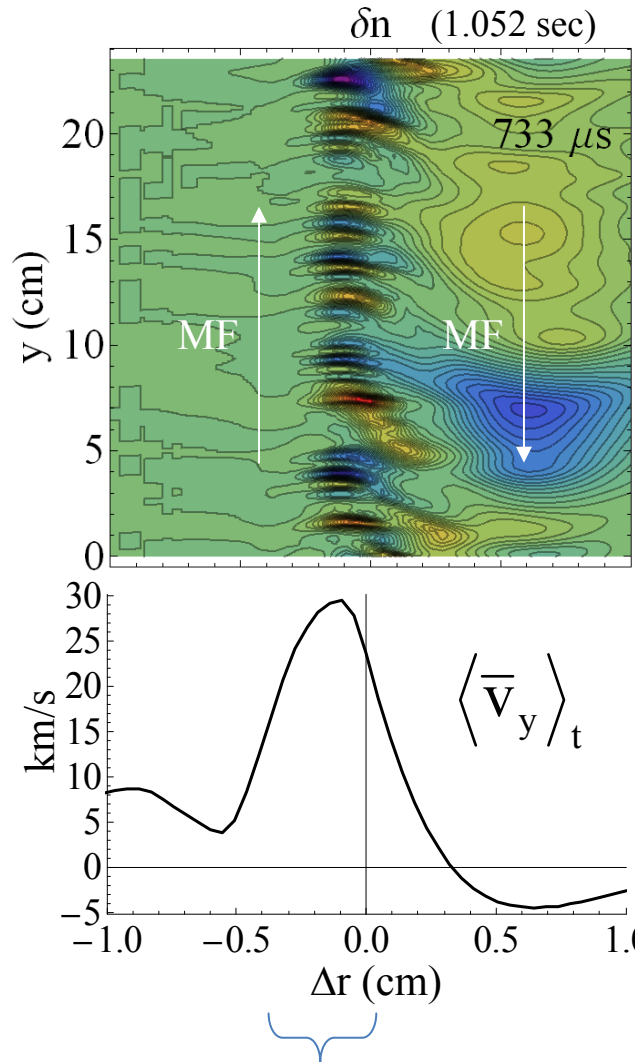
$$\langle q \rangle_{y,t} / (n_e c_s T_e)_{ref}$$



1.052 sec
SOLT

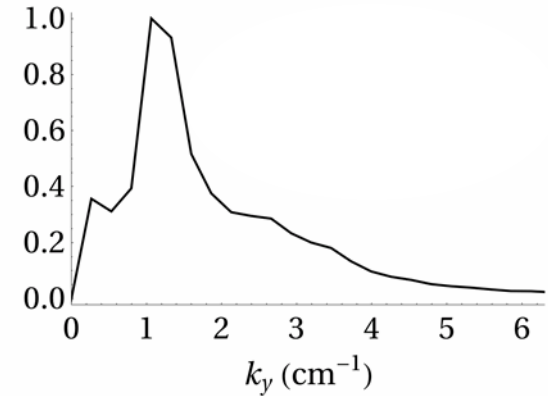
SOLT's QCM

The saturated turbulent state consists of a string of blobs, radially-localized about a maximum of the mean flow (MF) just inside the SEP, intermittently spilling plasma into the SOL where the flow reverses.



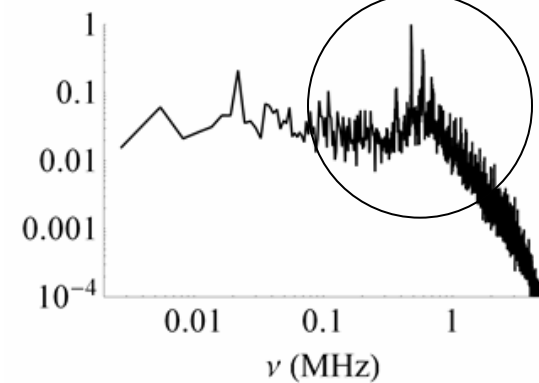
$$\langle |\delta n(k_y)|^2 \rangle_{t, \Delta r}$$

max @ $k_y = 1.06 \text{ cm}^{-1}$



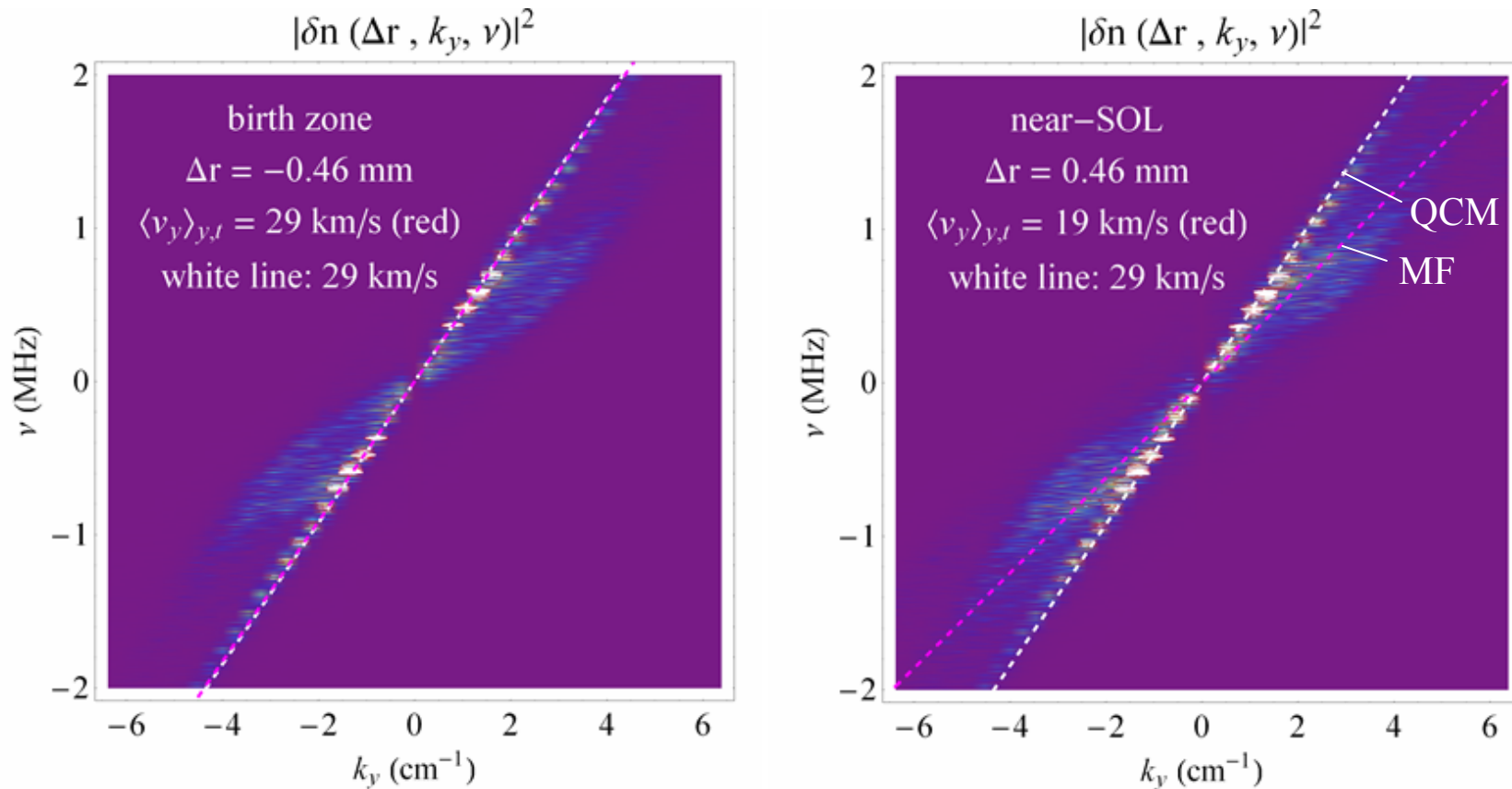
$$\langle |\delta n(\nu)|^2 \rangle_{y, \Delta r}$$

max @ $\nu = 479 \text{ kHz}$



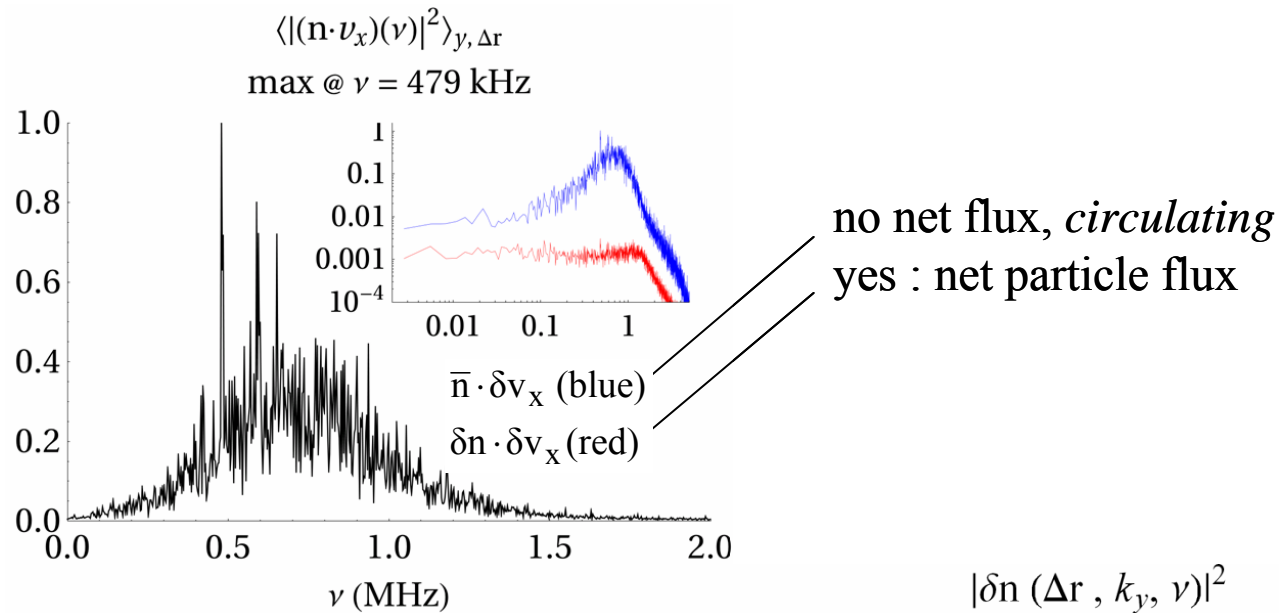
birth zone

QCM dispersion is established in the birth zone.



The local Doppler frequency corresponds to the QCM dispersion line (bright beads) only in the **birth zone**, where the time-averaged MF is maximized (flow shear = 0).

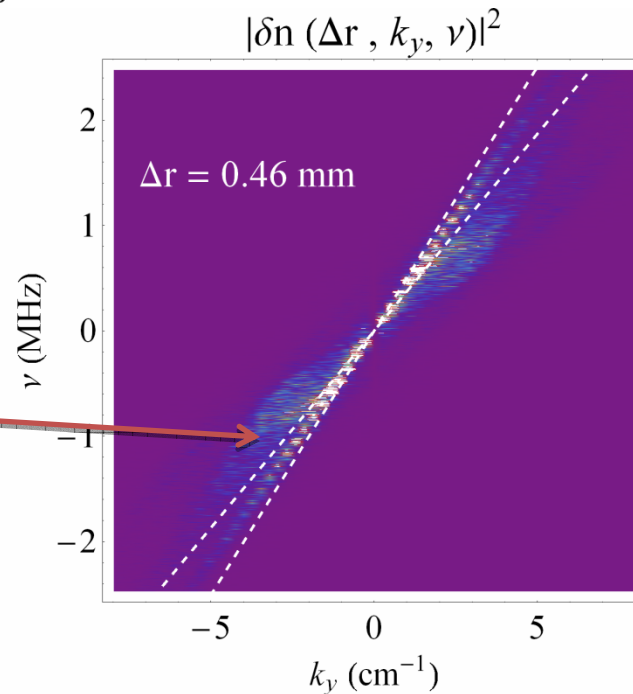
Particle Flux from the QCM



44% of the net particle flux,

$$\langle \Gamma \rangle_{y,t} = \sum_{k_y, \nu} \delta n(k_y, \nu) \cdot \delta v_x(k_y, \nu)^*,$$

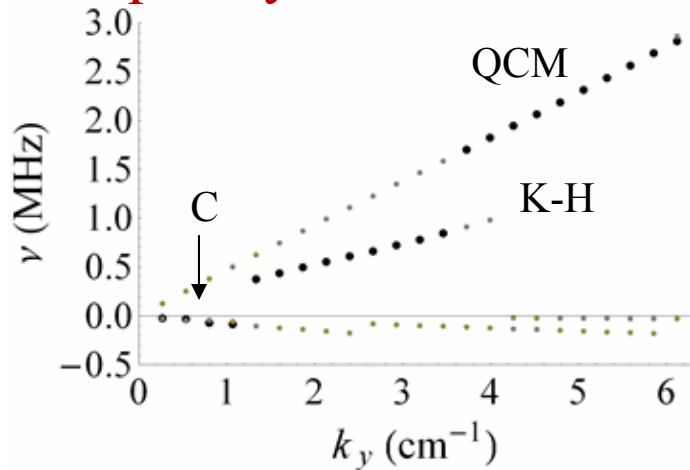
comes from inside the wedge
that includes the QCM.



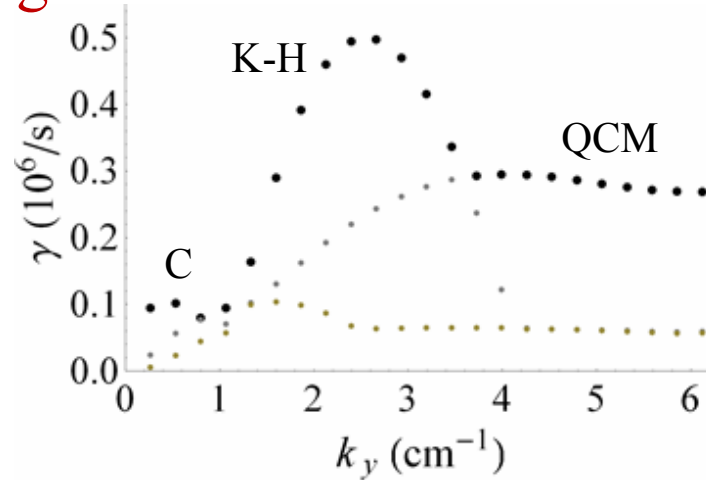
Linear Analysis of Time-Averaged Profiles

suggestive of underlying transport dynamics

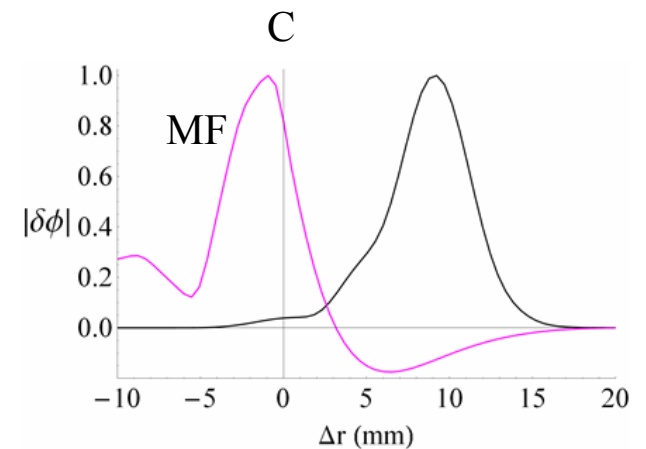
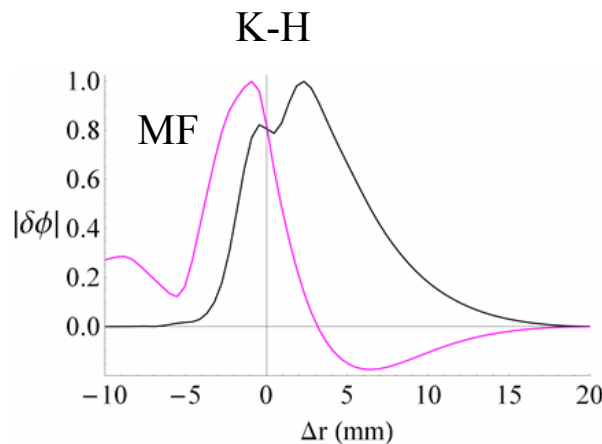
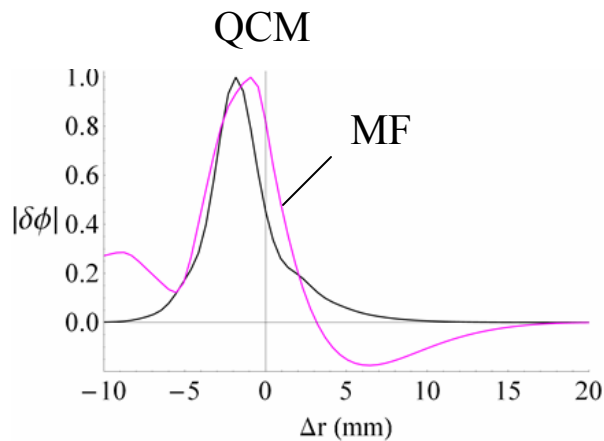
frequency



growth rate



radial eigenfunctions, $|\delta\phi(\Delta r)|$:



- drift-interchange mode
- blob birth zone

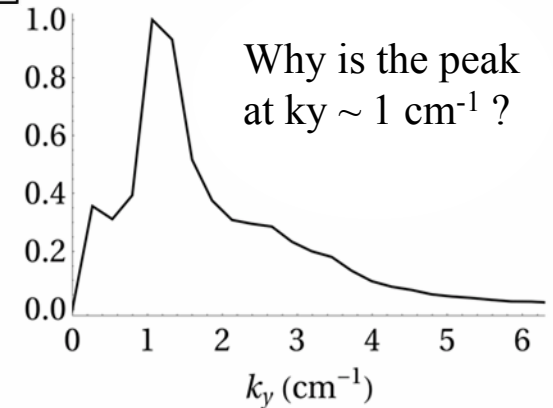
- straddles (\pm) MF regions
- blob emission

- sheath mode
- blob graveyard

Beyond Linear Analysis: The Vorticity Cascade Barrier

$$\langle |\delta n(k_y)|^2 \rangle_{t, \Delta r}$$

max @ $k_y = 1.06 \text{ cm}^{-1}$

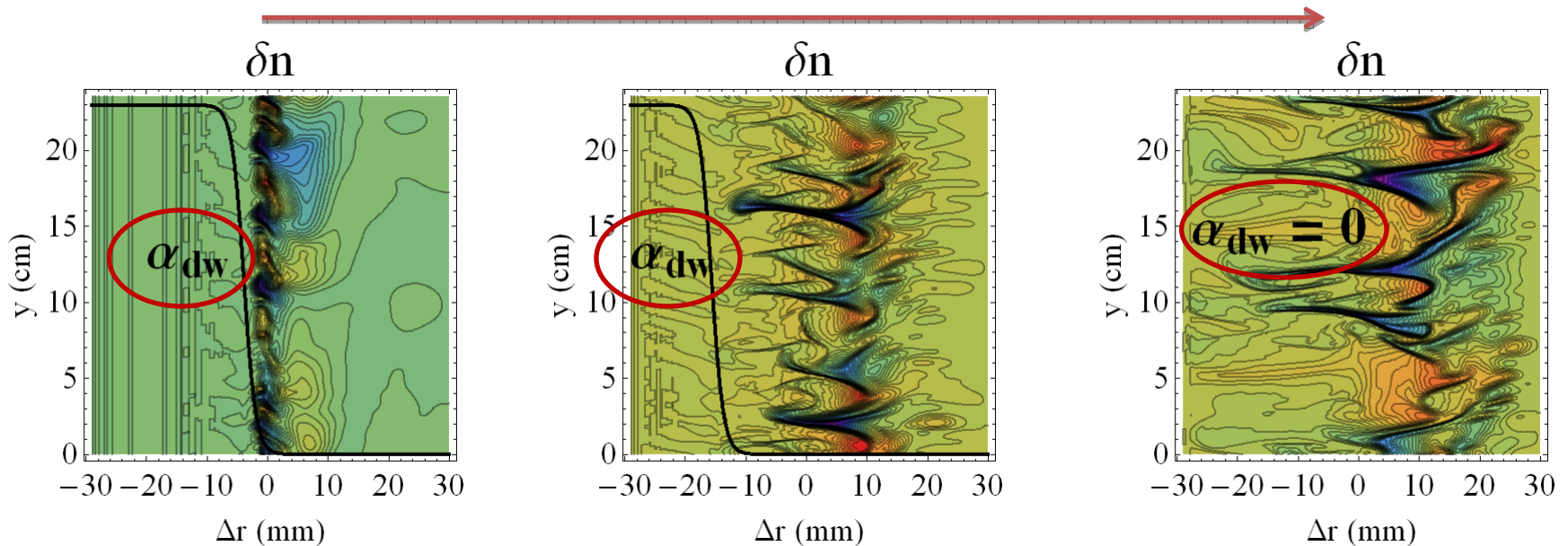


$$\frac{d}{dt} \nabla^2 \tilde{\Phi} = \left\{ \alpha_{dw}(\Delta r) \frac{\bar{T}^{3/2}}{\bar{n}} (\Phi - T \ln n) + \dots \right\}$$

α_{dw} : Drift Wave
Adiabaticity Parameter

- $\sim k_{//}^2 / v_{ei}$
- $f(\text{B field topology})$

Moving the DW region further into the core allows vorticity cascade to smaller k_y .



➤ This is consistent with the sharp change in the sign of the cross-phase ($\delta n, \delta \phi$) near the inflection point in the α_{dw} profile.

Summary

Part 1

Scaling of the SOL width for parallel heat flow

- Matching P_{SOL} with SOLT simulations, by adjusting the mean flow (τ)
 $\Rightarrow q_{//}$ - width scaling with T_e :
- $q_{//}$ is limited by collisions in the near-SOL: $q_{//, \text{CL}} \sim T_e^{7/2}$
 - consistent with T_e – dependence observed for this shot
 - differs from a similar study of NSTX scaling in the sheath-limited regime
 - (note: sheath-limited heat flux dominates in the far-SOL)

Part 2

Quasi-Coherent Mode

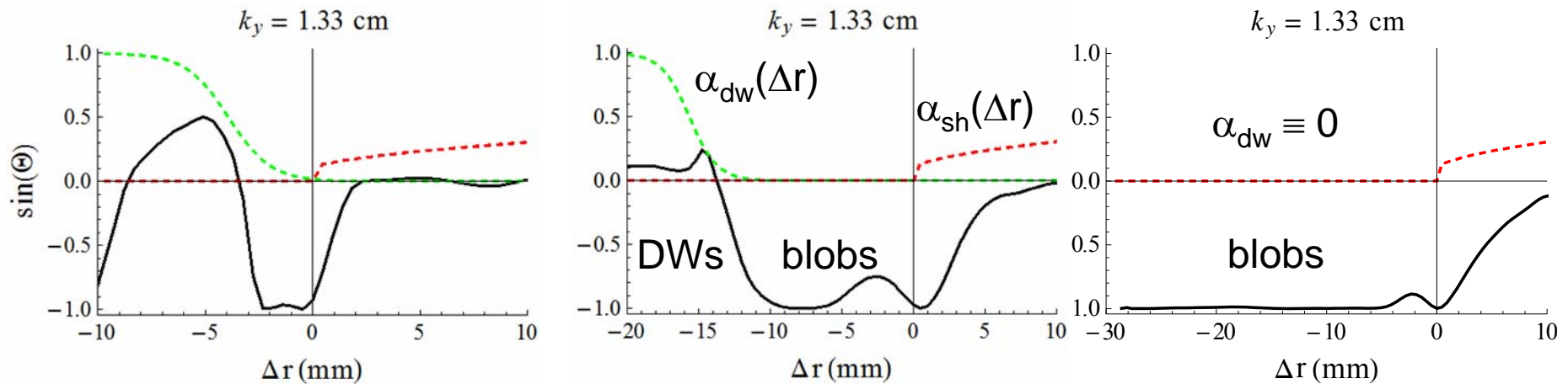
- A string of quasi-stationary blobs, moving with the mean flow in the edge
 - centered in the birth zone, where the mean flow shear rate = 0
 - energy spectrum consistent with experiment, $k_y \sim 1 \text{ cm}^{-1}$
 - accounts for 44% of the net particle flux, consistent with sustaining the EDA H-mode
 - linear unstable modes (drift-interchange, K-H) drive transport in the saturated state
 - drift-wave transition region is a barrier to vorticity cascade $\Rightarrow \langle |\delta n(\mathbf{k}_y)|^2 \rangle$ peaks at $k_y > 0$

extra
slide

The phase relation between $\delta\phi$ and δn changes abruptly at the entrance to the DW region.

Cross-Phase

$$\sin(\Theta) = \left\langle \delta n(\Delta r, k_y, t) \cdot \delta\phi(\Delta r, -k_y, t) \right\rangle_t / |(top)|$$



$\sin(\Theta) < 0$: conducive to the blob/hole generation and propagation paradigm
 $\sin(\Theta) > 0$: suppresses blob formation