
Sheared flow dynamics in edge turbulence*

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Introduction

- Core confinement in tokamaks is sensitive to edge profile gradients.
- The edge profiles are set by turbulence, so it is important to understand the mechanisms of turbulence saturation and how these influence the edge gradients.
- For typical edge profiles, there is no scale separation between background and turbulence, $k_{\perp}L_n \sim 1$.
- Edge turbulence studies require a “global” code (with edge and SOL profiles free to evolve), not a “local” or flux-tube code. Here, we use the 2D SOLT (SOL turbulence) code.
- We are studying the roles of
 - *pressure profile modification*
 - *Reynolds stress and sheared flows*
 - *radial variation in geometry and physics (open and closed field lines, drift wave, curvature driven and sheath regions)*

Outline

- *Quasi-linear calculations for Kelvin-Helmholtz (KH) and interchange modes*
- *Overview of SOLT turbulence simulations*
- *Comparison of quasi-linear theory with simulations*
- *Conclusions*

Quasi-linear model

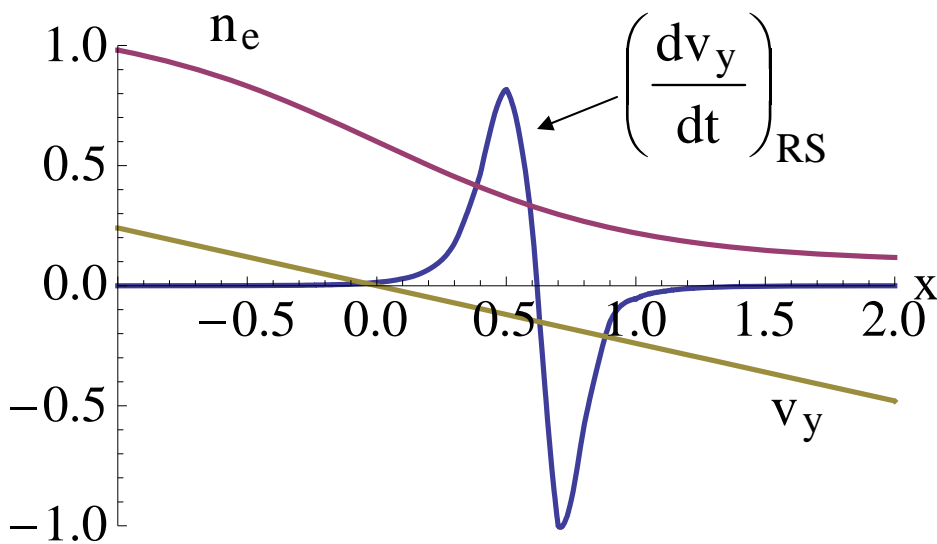
- **Interchange and Kelvin-Helmholtz (KH) stability is given by**

$$(\omega - k_y v_y)(\Phi'' - k_y^2 \Phi) + k_y v_y'' \Phi + \frac{k_y^2 n_0'}{n_0(\omega - k_y v_y)} \Phi = 0$$

- **BCs: $\Phi \rightarrow 0$ as $x \rightarrow \pm\infty$**
- **Inputs: k_y and profiles $n_0(x)$ and $v_y(x) \Rightarrow$ scalelengths L_n and L_v**
- **Compute the Reynolds' stress (RS)-generated $dv_y(x)/dt$.**
 \Rightarrow the *sign* of $dv_y(x)/dt$ depends on $k_y L_n$ and L_v/L_n .
- **Let $L_n = 1$ so that length scales are normalized to L_n .**

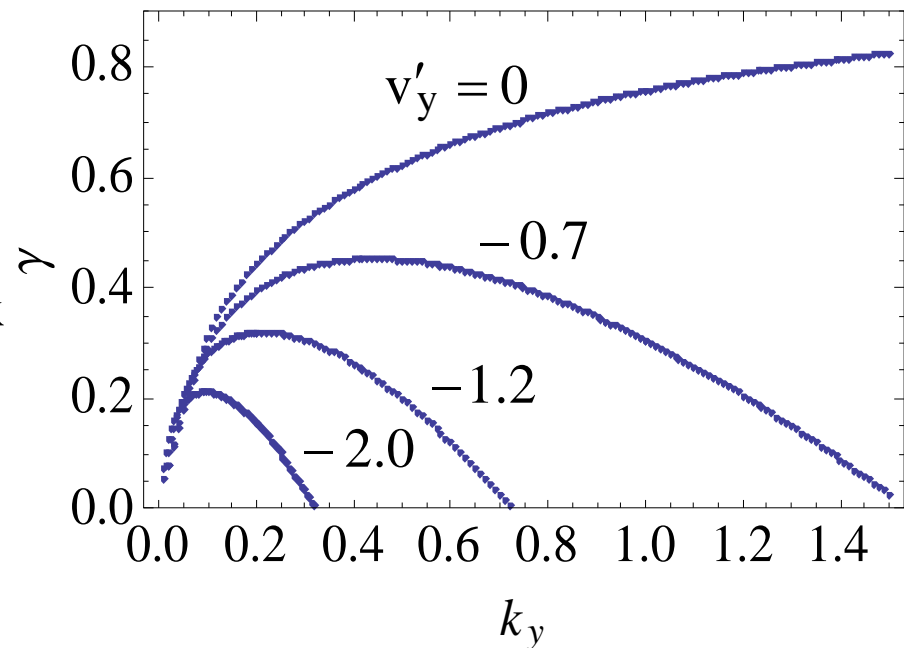
Constant velocity shear ($L_v \gg L_n$)

$k_y L_n = 0.6$



RS tries to increase the v_y gradient at the mode center \Rightarrow ZF instability

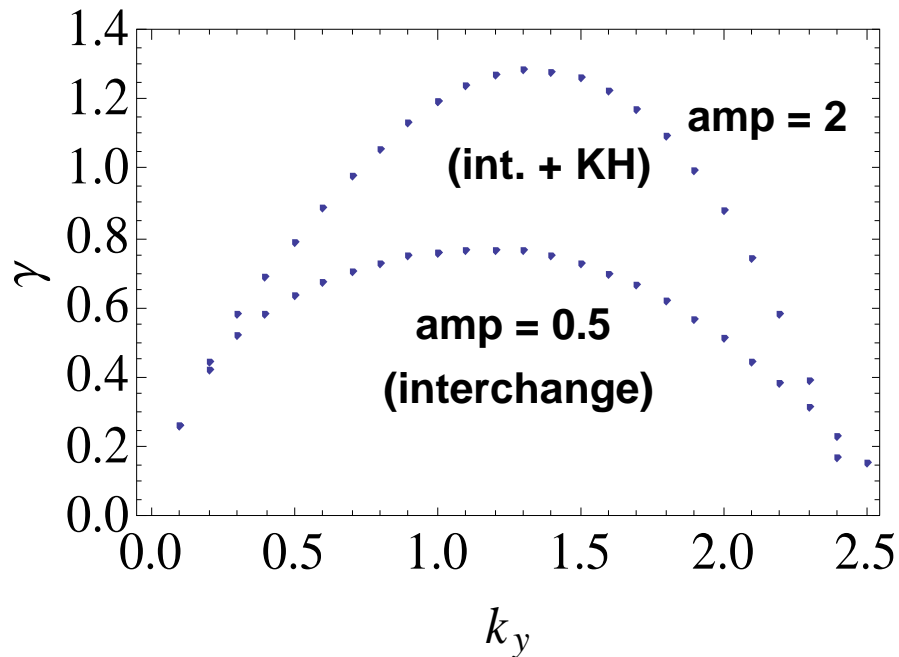
Interchange growth rate γ vs $k_y L_n$



Constant v_y shear \Rightarrow no KH mode.

Shear stabilizes the interchange mode and is most effective at high k_y .

Localized v_y shear ($L_v \sim L_n$)

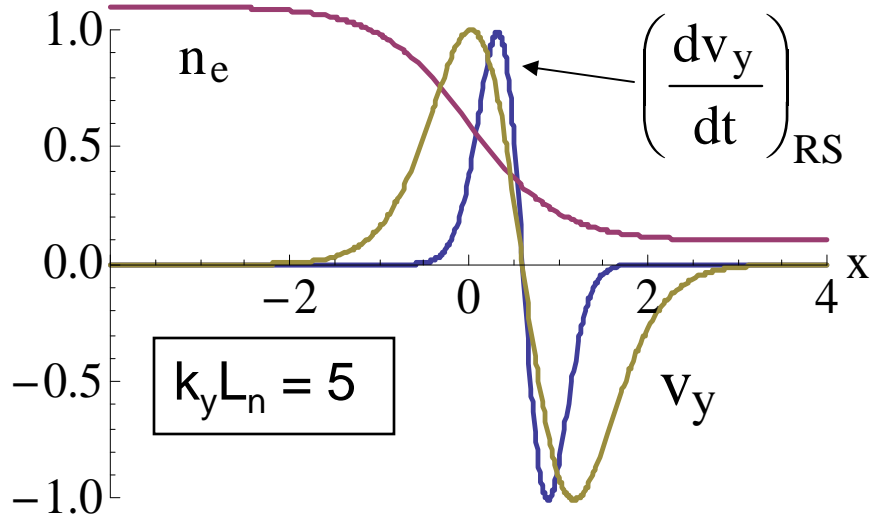


The amp = 0.5 case is largely interchange driven.

Note dependence of growth rate on $k_y L_n$:

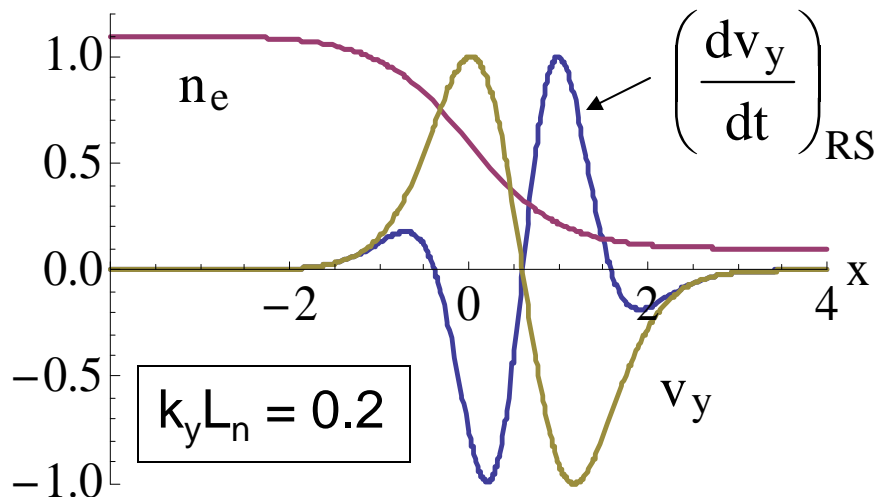
- (a) small $k_y L_n \Rightarrow \gamma \rightarrow 0$ for all instability drives
(flows are not important)
- (b) intermediate $k_y L_n \Rightarrow$ flows are destabilizing (KH mode)
- (c) large $k_y L_n \Rightarrow$ flows stabilize interchange mode

Effect of the RS depends on $k_y L_n$



amp $\ll 1$ for both cases \Rightarrow **pure interchange instability drive**

$k_y L_n \gg 1 \Rightarrow$ RS tries to increase the gradient in v_y at the mode center, enhancing the flow that was present.



$k_y L_n \leq 1 \Rightarrow$ The RS opposes the existing flow.

Quasi-linear picture

From these and other cases, varying the shapes of the profiles and the scale lengths, we obtain the following picture of how curvature-driven interchange modes interact with sheared flows:

- **Pure KH modes** act to reduce the v_y shear driving them (at least locally, although this may increase the shear elsewhere).
- Shear in v_y stabilizes the interchange mode and is most effective on the **high k_y** modes.
- The Reynolds Stress (RS) from **pure interchange modes** acts to enhance the imposed flows if the flows are **large-scale** with respect to the mode's radial variation, while the RS opposes the seed flows if those flows are **small scale**.

Quasi-linear picture - 2

- An important implication of this work is that **long wavelength modes (with respect to the pressure gradient) cannot be stabilized by self-generated sheared flows; instead, they must be stabilized by pressure profile modification (wave-breaking, plateau formation)**. This agrees with the SOLT simulations described subsequently.
- Finally note that these results suggest that the inverse cascade in k_y may be responsible for this behavior:
 - **Large k_y** , associated with small-scale radial mode structure, **feeds** the (larger-scale) flow.
 - **Small k_y** , associated with large-scale radial mode structure, **saps** the (smaller-scale) flow.
- This picture is in qualitative agreement with SOLT simulations of turbulence saturation as a function of radial gradients

Quasi-linear picture - 3

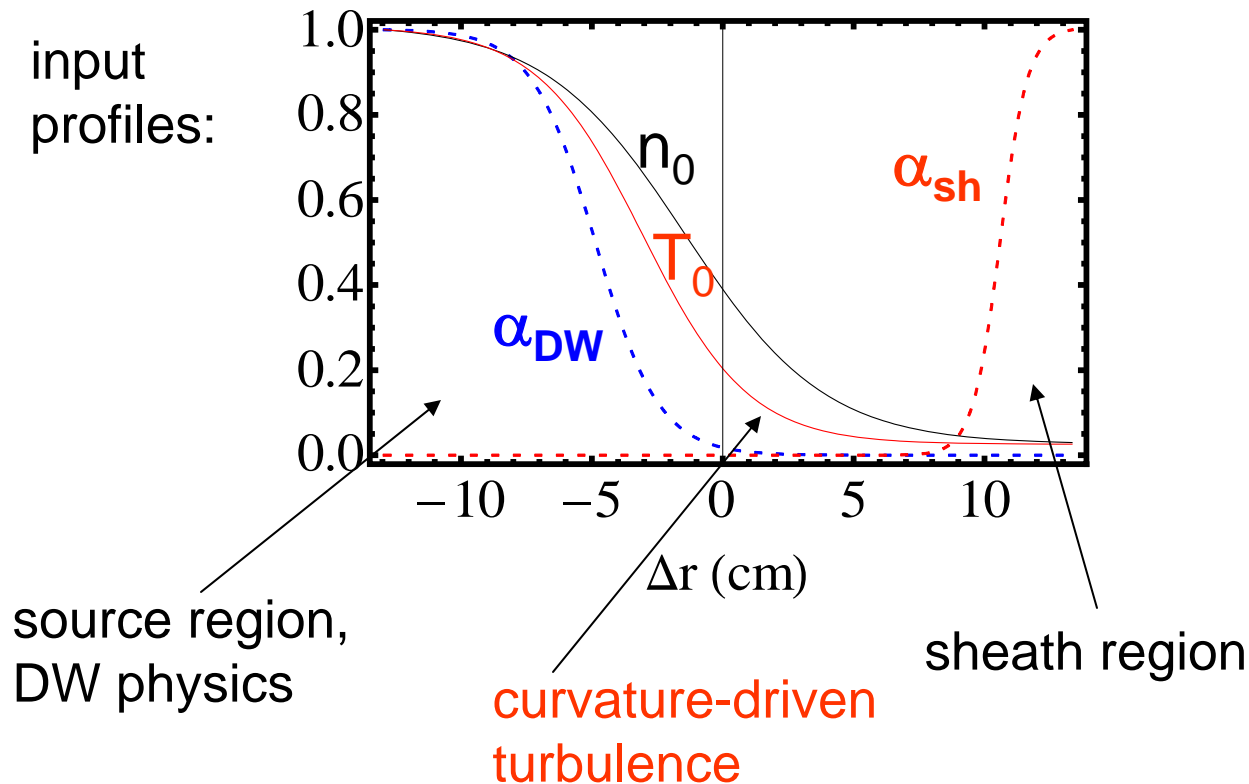
For comparison with the SOLT turbulence simulations, summarize the results of the quasi-linear calculations for **curvature-driven interchange modes** in the following table:

	$L_v \leq L_p$	$L_v \gg L_p$
$k_y L_p \gg 1$	RS+	RS+
$k_y L_p \leq 1$	RS-	RS+

where RS+ (RS-) \Rightarrow the Reynolds Stress (RS) acts to **increase** (**decrease**) the flow shear at the mode center.

RS+ is the usual ZF instability found in flux-tube (local) simulations.
RS- requires a global turbulence code (e.g. SOLT).

SOLT turbulence studies



Definitions:

$$\alpha_{dw} = \frac{2\rho_s m_i c_s}{L_{\parallel e}^2 v_{ei0} m_e}$$

$$\beta = \frac{2\rho_s}{R} \quad \alpha_{sh} = \frac{2\rho_s}{L_{\parallel s}}$$

$$\gamma_0 \equiv \left(\frac{\beta}{n} \frac{dP}{dx} \right)^{1/2}$$

We have done a study of the saturation of the turbulence as a function of the instability drive $\gamma_0 \propto dP/dx$ ($P = n T$)

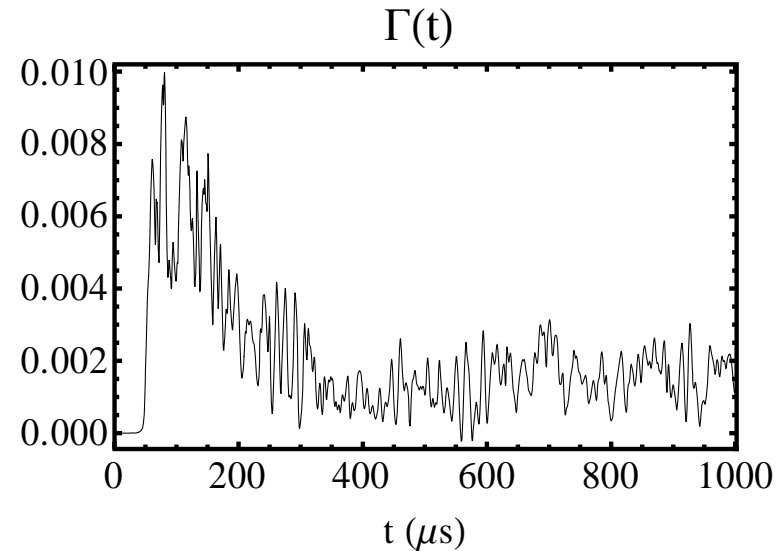
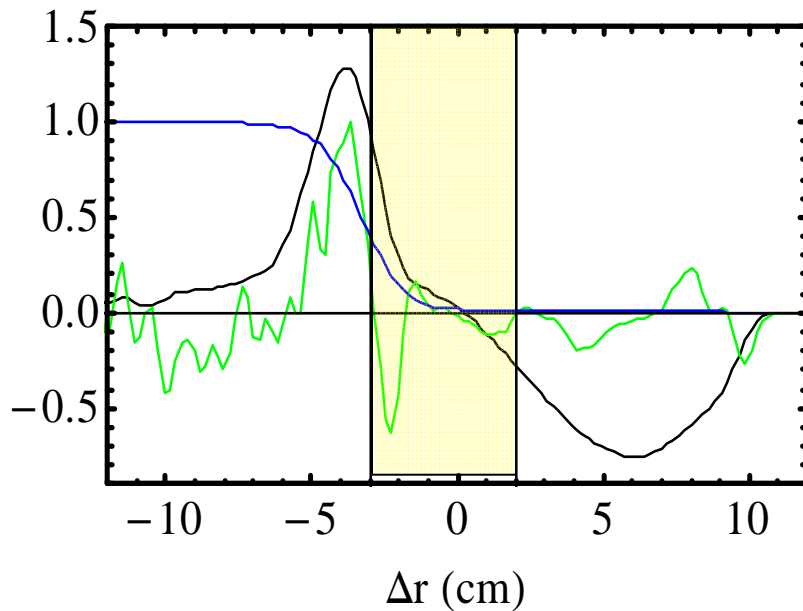
Large gradient case

run32

Quasi-steady state:

Below: shaded area indicates FWHM of particle flux $\Gamma = \langle n v_x \rangle$

\bar{p} (blue), \bar{v}_y (black), $\overline{dp_y/dt}$ (green)



⇒ **RS- regime** near max Γ

⇒ flows are not strong enough to stabilize interchange modes)

$$\langle v'_y \rangle_{y,t} / \langle \gamma_{\text{mhd}} \rangle_{y,t} \approx 0.6$$

$$\langle k_y L_p \rangle_{y,t} \approx 1.2$$

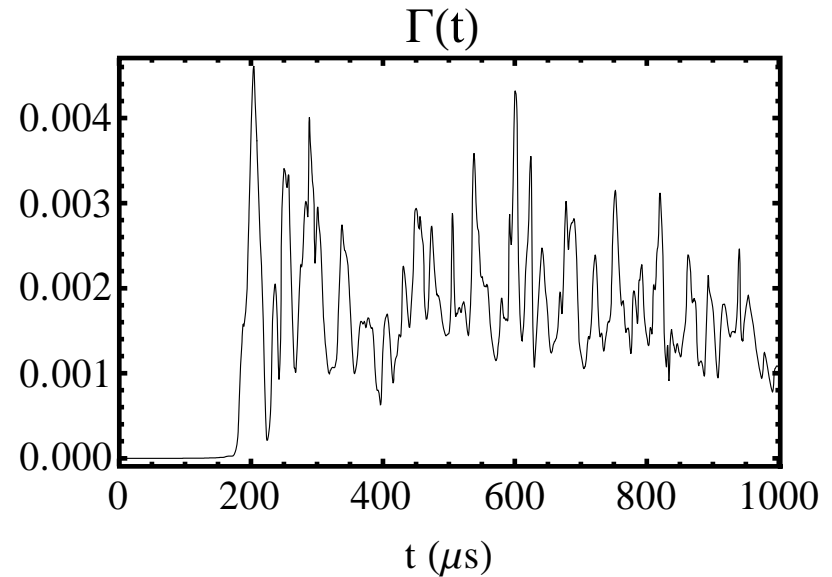
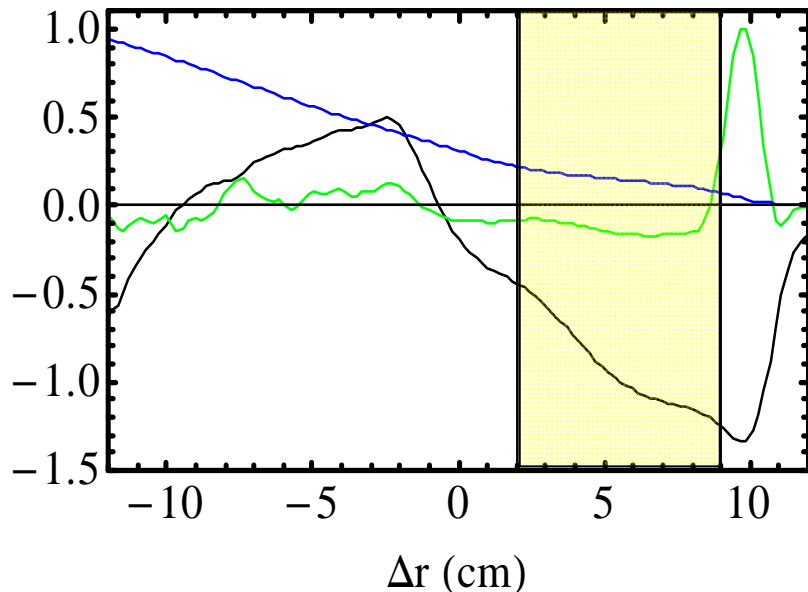
Small gradient case

run33

Right: note different initial transient and fluctuation behavior than in run32

Below: shaded area indicates FWHM of particle flux $\Gamma = \langle n v_x \rangle$

\bar{p} (blue), \bar{v}_y (black), \overline{dp}_y/dt (green)



⇒ **RS+ regime** near max Γ

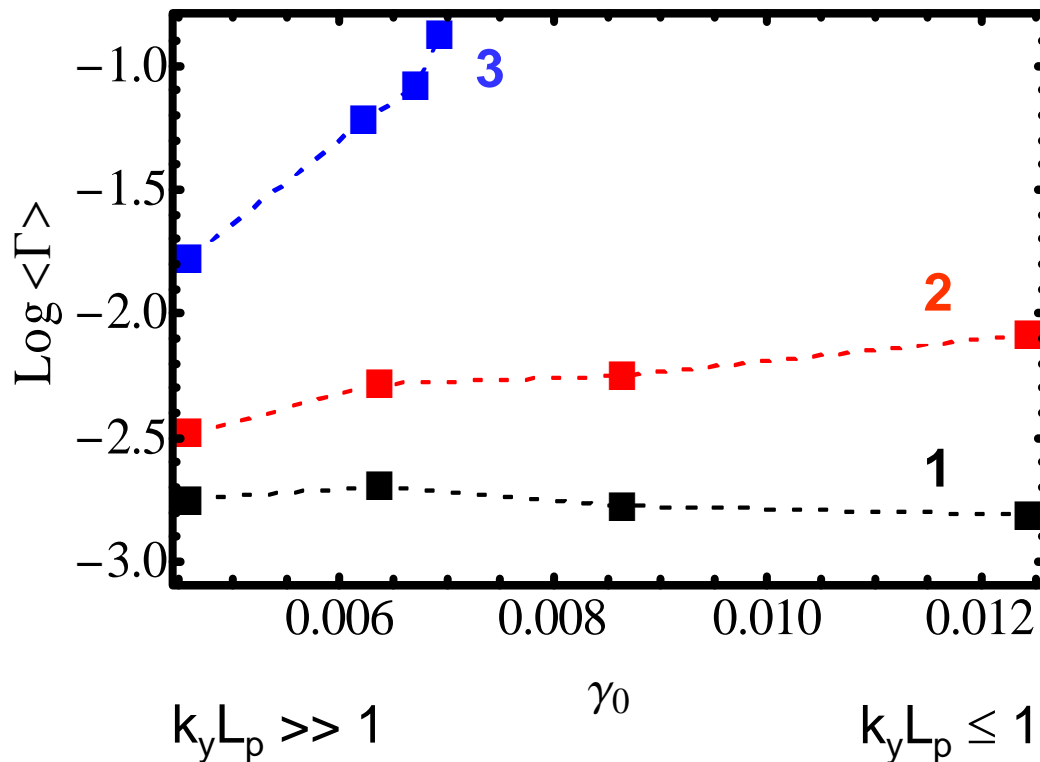
⇒ **sheared flows help stabilize interchange modes**

$$\langle v'_y \rangle_{y,t} / \langle \gamma_{\text{mhd}} \rangle_{y,t} \approx 2.4$$

$$\langle k_y L_P \rangle_{y,t} \approx 7.5$$

Regime diagram for SOLT study

Plot of particle flux Γ vs interchange growth rate γ_0 (see p. 11)



Saturation mechanisms

1 \Rightarrow flow + profile mod

2 \Rightarrow profile mod only

3 \Rightarrow flow only (frozen profiles)

- y- and t-averaged flux
- evaluated at radial location of max flux Γ

blue curve \Rightarrow stabilizing effect of sheared flow decreases rapidly as grad-P increases and $k_y L_p$ decreases (agrees with the quasi-linear model)

Conclusions: emerging picture

- Here we have studied **curvature-driven interchange turbulence** in the edge plasma, saturated by a combination of
 - sheared poloidal flows
 - radial profile flattening
- The dominant saturation mechanism depends on the radial gradients.
- The experimental geometry (changing topology of magnetic field lines) sets the natural radial scale L_x of the edge turbulence.
- The radial mode structure of the turbulence, and profiles, determine the k_y spectrum, and $k_y L_x$ determines the evolution of the sheared flows.
- **The relative ordering of k_y , L_p , and L_v (pressure and sheared velocity scale lengths) influences the sign of the Reynolds Stress (RS) and whether it causes the sheared flows to grow (RS+) or damp (RS-)**
 - this may be related to the inverse cascade
 - see pp. 8 – 10 for details

Conclusions (cont.)

- The **RS regime influences the saturation mechanism** for the interchange turbulence
 - RS+ \Rightarrow shear flow stabilization
 - RS- \Rightarrow profile flattening
- When the mode significantly overlaps with the DW region inside the edge or the sheath-connected region in the far SOL, **additional physics** enters:
 - DW \Rightarrow directionality, poloidal flows
 - sheath \Rightarrow sink for particles, heat and momentum;
Bohm sheath-potential-driven flows
- Preliminary **comparisons with SOLT turbulence simulations** are encouraging but more work is needed to obtain optimal simulations for this comparison.